

TlGa_{1-x}Er_xSe₂ SOLID SOLUTIONS, THEIR ELECTRICAL AND OPTICAL PROPERTIES

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The complex permittivity, dielectric loss tangent, and conductivity σ_{ac} of TlGa_{1-x}Er_xSe₂ solid solutions ($x = 0; 0.001; 0.005$) were studied in the temperature range 150-300 K in an alternating electric field (25-10⁶ Hz). The change in dielectric constant as a function of temperature in the studied crystals is due to the presence of low-frequency relaxation polarization. Conductivity obeys the law $\sigma_{ac} \sim f^{0.8}$ at frequencies of 10²-10⁴ Hz, which, in the studied crystals, indicates the presence of hopping conductivity by states localized near the Fermi level. The parameters of the states localized in the TlGa_{1-x}Er_xSe₂ band gap are estimated. It was shown that an increase in the Er impurity leads to an increase in the density of localized states in TlGaSe₂ and a decrease in their energy spread. The optical absorption edge of TlGa_{1-x}Er_xSe₂ crystals was studied, and the temperature dependence of the band gap in them was obtained.

Keywords: solid solutions, complex dielectric permittivity, frequency dispersion, dielectric loss, conductivity, optical absorption edge.

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1. INTRODUCTION

TlGaSe₂ crystal has a layered structure of A^{III}B^{III}C^{VI}₂ thallium chalcogenide type compounds. This crystal has a monoclinic structure, C2/c-C_{2h}⁶ space group, with unit cell parameters $a = 10.772$, $b = 10.771$, $c = 15.636$, $\beta = 100.6^\circ$ [1].

TlGaSe₂ is a p-type wide-gap semiconductor and possesses practically important physical characteristics: high photo- and X-ray sensitivity. On its basis, photoelectric converters, spectrum analyzers, X-ray, gamma, and neutron detectors [2,3,4] are developed, and the possibility of creating a phototransistor is shown [5].

In [6,7], hopping conductivity in TlGaSe₂ single crystals was studied. The results of studying the effect of γ radiation on the conductivity and dielectric characteristics of TlGaSe₂ single crystals are presented in [8,9]. Many works, for example [10,11], are devoted to the study of the optical properties of the TlGaSe₂ compound.

Specific feature of A^{III}B^{III}C^{VI}₂ family crystals is that these semiconductors also possess ferroelectric properties. Due to the layered structure, polytypic modifications occur, resulting in changes in their physical properties [12]. In addition, a large number of crystals of this type are characterized by consecutive phase transitions [13,14].

Our assessment of the solubility of erbium in the TlGaSe₂ lattice, taking into account the known effective ionic atomic radii, indicates that the radius of the implanted Er³⁺ impurity atom (1.03Å) is closer to the Ga³⁺ radius (0.76Å) than to the Tl¹⁺ radius (1.64Å), i.e. the partial replacement of gallium with erbium in layered TlGaSe₂ crystals corresponds to the condition for the formation of a substitution solution. The aim of this work is to study the electrical, dielectric, and optical properties of TlGa_{1-x}Er_xSe₂ solid solutions ($x = 0; 0.001; 0.005$).

2. METHODS OF THE INVESTIGATION

The synthesis of TlGa_{1-x}Er_xSe₂ solid solutions ($x = 0; 0.001; 0.005$) was carried out by fusing highly pure chemical elements taken in stoichiometric ratios in quartz ampoules evacuated to 10⁻³ Pa. Samples were homogenized by annealing at a temperature of 750K for 120h. Single crystals of TlGa_{1-x}Er_xSe₂ were grown from the synthesized compounds by the Bridgman method [15]. The crystals were layered and easily chipped along the basic plane. X-ray phase analysis of the obtained samples was carried out on a D8-ADVANCE diffractometer in the mode 0.5° < 2 θ < 80° (CuK α radiation, $\lambda = 1.5418\text{Å}$).

The dielectric constant ϵ and the dielectric loss tangent $\tan\delta$ were measured using the capacitor method in the temperature range 150-300 K using an E7-25 digital device on single-crystal wafers 0.2-0.3 mm thick. The contacts were made of silver paste. The frequency range of the alternating electric field was 25-10⁶ Hz.

The measurements of the capacity of the samples were carried out with an accuracy of ± 0.01 pF, and the error in measuring the quality factor ($Q = 1/\tan\delta$) was 0.001%.

The electrical conductivity σ was measured along the layers by the four-probe method in the mode of continuous increase (decrease) in the sample temperature at a rate of ≈ 0.1 K/min. The amplitude of the alternating electric field applied to the samples was within the ohmic region of the current-voltage characteristic.

Samples for studying the optical absorption spectra of layered semiconductor crystals of TlGa_{1-x}Er_xSe₂ ($x = 0; 0.001; 0.005$) were chipped from a single-crystal ingot and had the form of thin plates with a thickness of 20 to 140 μm . Light was directed to the samples parallel to the crystallographic axis c . The optical transmission spectra were studied using an apparatus based on an MDR-23 monochromator and a

nitrogen cryostat with temperature stabilization in the range 77–300K (stabilization accuracy ±0.5K). The resolution of the installation was 2Å.

To measure the optical absorption coefficient in a wide range from 10 to 10⁴ cm⁻¹, we used the data of measurements of the intensity of the light beam passing through samples with different thicknesses, and to cover the entire interval, we had to break it into 3 sections and take into account the transmission measurements of three pairs of samples of corresponding thicknesses for each composition studied. For each plot, the absorption coefficient was calculated by the formula $\alpha = 1/(d_2 - d_1)\ln(I_1/I_2)$, where d_1 and d_2 are the thicknesses of the samples, and I_1 and I_2 are the intensities of the light transmitted through them. Since the value of αd was greater than unity for each sample and the corresponding section, the interference of light beams passing through and reflected from the back plane of the crystal plate was very weak, and we did not observe it.

3. EXPERIMENT RESULTS AND DISCUSSION

When processing X-ray data of TiGa_{1-x}Er_xSe₂ solid solutions, it was found that the amount of Er impurity element used ($x = 0.001; 0.005$) has little effect on the lattice parameters of the initial TiGaSe₂ compound. However, the addition of Er to TiGaSe₂ leads to a change in the color of the initial crystal to a darker one (see Fig.1).

The results of X-ray analysis of the crystals under study are presented in Tables 1, 2 and 3. X-ray studies showed that TiGa_{1-x}Er_xSe₂ ($x = 0; 0.001; 0.005$) single crystals have a monoclinic structure, space group C_s^4 , and the following lattice parameters were obtained: $a = 10.744 \text{ \AA}$, $b = 10.773 \text{ \AA}$, $c = 15.623 \text{ \AA}$, $\beta = 100.04^\circ$, $z=16$, which is in good agreement with [1]. The frequency dependences of the conductivity of TiGa_{1-x}Er_xSe₂ crystals at room temperature are shown in the Fig.2.



Fig. 1. Single crystals grown by the Bridgman method: TiGaSe₂ on the left, TiGa_{0.995}Er_{0.005}Se₂ on the right.

Table 1.

Calculation of the X-ray diffraction pattern of the TiGa_{0.999}Er_{0.001}Se₂ crystal.

N ₀	2θ	I/I ₀ , %	d _{exp} (Å)	d _{calc} (Å)	hkl
1	11°46'	17,4	7,715	7,695	200
2	23°11'	100	3,845	3,878	422
3	23°6'	8,1	3,767	3,778	122
4	24°95'	2,8	3,566	2,568	222
5	26°30'	2	3,386	2,398	031
6	27°54'	6,4	3,236	2,236	322
7	31°09'	18,9	2,875	2,878	422
8	33°82'	2,9	2,648	2,650	140
9	34°97'	31,8	2,564	2,544	522
10	37°80'	2,1	2,378	2,382	404
11	47°26'	7,7	1,922	1,924	215
12	47°80'	5,8	1,901	1,902	144
13	53°84'	2,8	1,702	1,703	740
14	56°27'	2	1,634	1,634	831
15	58°13'	2,1	1,586	1,587	362

Table 2.

Calculation of the X-ray diffraction pattern of the $TlGa_{0,995}Er_{0,005}Se_2$ crystal.

№	2θ	I/I ₀ , %	d _{exp} (Å)	d _{calc} (Å)	hkl
1	11°44'	13,8	7,726	7,625	200
2	23°11'	100	3,846	3,847	400
3	23°52'	8,3	3,779	3,777	122
4	24°98'	2,4	3,562	2,568	222
5	26°33'	2,5	3,382	2,392	031
6	27°54'	6,4	3,237	2,237	322
7	31°05'	18,5	2,878	2,878	422
8	33°83'	4	2,648	2,650	140
9	35°01'	43,1	2,561	2,565	600
10	35°40'	2,5	2,534	2,544	522
11	41°42'	1,8	2,178	2,185	711
12	47°25'	7,9	1,922	1,923	802
13	47°81'	5	1,901	1,902	144
14	50°52'	1,6	1,805	1,806	153
15	53°88'	3,4	1,700	1,702	740
16	56°18'	1,9	1,636	1,634	426
17	58°15'	1,9	1,585	1,587	362
18	64°14'	1,4	1,451	1,437	835
19	65°05'	1,8	1,433	1,433	264

Table 3.

Crystal lattice parameters of solid solutions of the $TlGaSe_2$ - $TlErSe_2$ system.

Compound	Crystal lattice system	Crystal lattice parameters, (Å)			The number of formula units in a unit cell, Z
		a	b	c	
$TlGa_{0,999}Er_{0,001}Se_2$	monoclinic	10,744	10,773	15,623	16
$TlGa_{0,995}Er_{0,005}Se_2$	monoclinic	10,744	10,773	15,623	16
$TlGaSe_2$	monoclinic	10,715	10,694	15,690	16

Table 4.

The parameters of localized states in $TlGa_{1-x}Er_xSe_2$ solid solutions obtained from high-frequency dielectric measurements.

Crystal composition $TlGa_{1-x}Er_xSe_2$	T, K	$N_F, eV^{-1}cm^{-3}$	$\Delta E, meV$	N_t, cm^{-3}
x = 0	180	$1.54 \cdot 10^{18}$	10	$1.54 \cdot 10^{16}$
	230	$1.76 \cdot 10^{18}$	9	$1.56 \cdot 10^{16}$
	300	$2.23 \cdot 10^{18}$	7	$1.57 \cdot 10^{16}$
x = 0.005	180	$2.48 \cdot 10^{18}$	6.7	$1.66 \cdot 10^{16}$
	230	$2.58 \cdot 10^{18}$	6.3	$1.62 \cdot 10^{16}$
	300	$3.08 \cdot 10^{18}$	5	$1.54 \cdot 10^{16}$

From this it can be seen that for compositions with 0.5% and 1% Er, the difference is small. Therefore, we studied the temperature dependences of the dielectric characteristics for the composition $TlGa_{0,995}Er_{0,005}Se_2$.

Fig. 3(a) ($TlGaSe_2$) and 3(b) (Er 0.5%) show the temperature dependence of the dielectric permittivity ϵ . It can be seen that, with increasing temperature, ϵ has a tendency to increase, especially at low frequencies.

Fig. 4 shows the temperature dependence of the electrical conductivity σ . For the $TlGa_{0,995}Er_{0,005}Se_2$ samples, as shown in the dielectric permittivity, the value of σ increases as the temperature increases. The absolute value of σ depends on the frequency of the measurement area, which is more pronounced at lower temperatures. The value of σ increases several times as the frequency increases. At low temperatures,

conductivity is practically independent of temperature. In the high temperature region, the conductivity increases exponentially with the law of $\sigma \sim e^{1/T}$. This indicates that at high temperatures σ mainly depends on the concentration of the main carriers.

Fig. 5(a) and 5(b) show the temperature dependence of the dielectric loss tangent ($\tan\delta$). As can be seen from the pictures, the values of the $\tan\delta$ of the TlGa_{0.995}Er_{0.005}Se₂ crystal increase as the temperature rises compared with TlGaSe₂. This increase can be related to the increased concentration of free charge carriers.

As can be seen from Fig.6, the value of the dielectric permittivity of TlGa_{1-x}Er_xSe₂ solid solutions decreases with increasing frequency. Such behavior of the dielectric permittivity in the studied crystals is due to the presence of low relaxation polarization at low frequencies, causing dielectric loss at low frequencies. In the $\tan\delta(T)$ curves, the semiconductor maximums

are visible and spread when the frequency of the measurement increases from 1kHz to 1MHz.

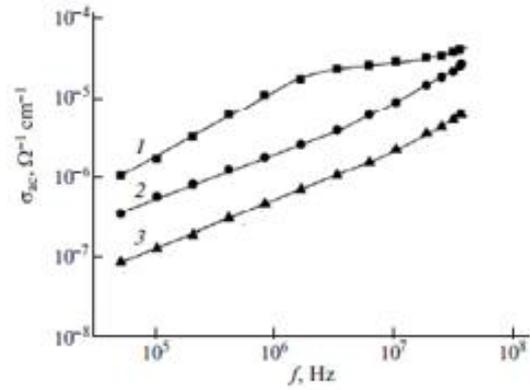


Fig. 2. Frequency dependent ac-conductivity of the TlGa_{1-x}Er_xSe₂ crystals: $x = 0$ (1), 0.001 (2), 0.005 (3). $T = 300$ K.

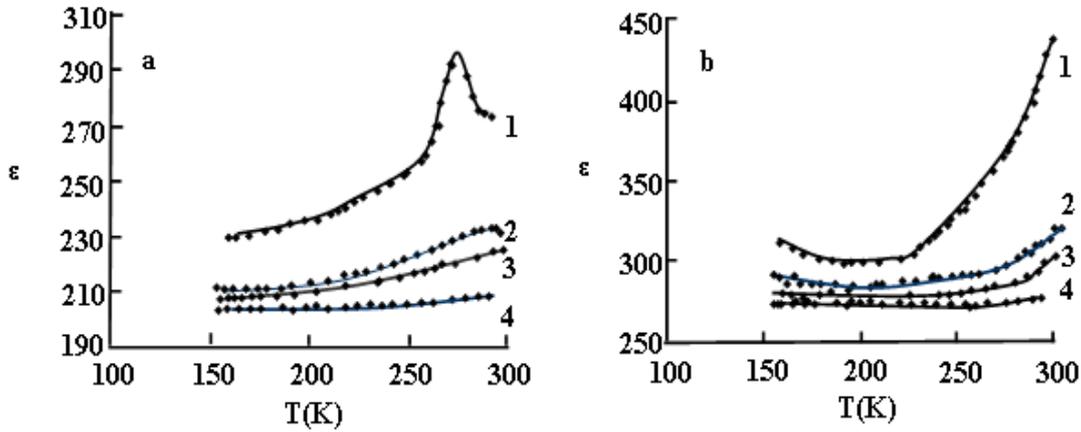


Fig. 3. Temperature dependence of dielectric permittivity of the TlGa_{1-x}Er_xSe₂ single crystal: a) $x=0$; b) $x=0.005$. (1- 1 kHz, 2- 10 kHz, 3- 100 kHz, 4- 1 MHz).

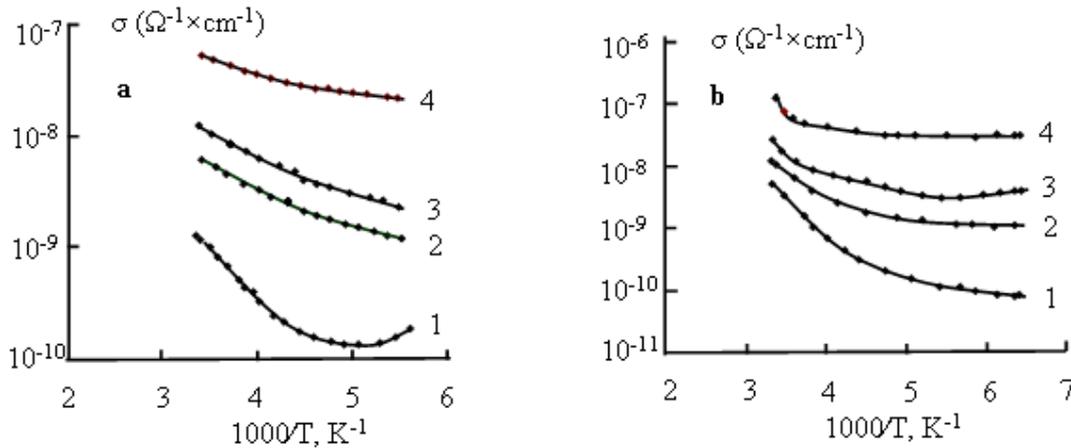


Fig. 4. Temperature dependence of the electrical conductivity of the TlGa_{1-x}Er_xSe₂ crystals: a) $x=0$; b) $x=0.005$. (1-1kHz, 2-10 kHz, 3-100kHz, 4-1MHz).

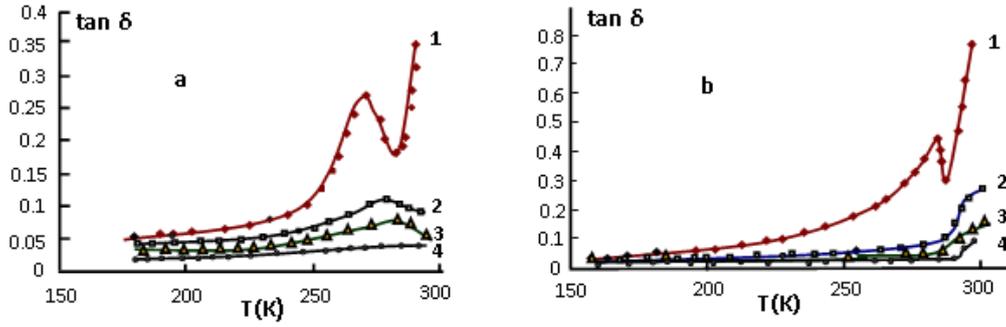


Fig. 5. Temperature dependence of the loss tangent in the $TiGa_{1-x}Er_xSe_2$ crystals: a) $x=0$; b) $x=0.005$. (1-1kHz, 2-10 kHz, 3-100kHz, 4-1MHz).

That is, these maximums are located in the temperature region, where the value of the dielectric permittivity is sharply increased. This anomaly is due to the ferroelectric phase transition in $TiGaSe_2$, which is also due to the dimensionless phase in the temperature region [16]. The real part of dielectric permittivity (ϵ') decreases with increasing frequency for $TiGaSe_2$ and $TiGa_{0.995}Er_{0.005}Se_2$. The increase in temperature, as well as the impurity of 0.5% Er in the $TiGaSe_2$ crystal, significantly increases the value of the ϵ' .

The frequency dependence of the imaginary part ϵ'' of the complex dielectric permittivity of $TiGa_{1-x}Er_xSe_2$ ($x=0; 0.005$) crystals at different temperatures is

shown in Fig. 7. It is determined that the frequency dependence of the imaginary part of the complex dielectric permittivity is of a relaxation nature.

Fig. 8 shows the comparison of the frequency dependence of the $\tan\delta$ of the $TiGaSe_2$ (curve 1) and the $TiGa_{0.995}Er_{0.005}Se_2$ (curve 2) crystals. This dependence is characterized by hyperbolic reduction, which is associated with a loss of conductivity. The increase of $\tan\delta$ at high frequencies indicates the presence of relaxation losses. Compared to $TiGaSe_2$, the value of the $\tan\delta$ of $TiGa_{0.995}Er_{0.005}Se_2$ crystal is significantly increased.

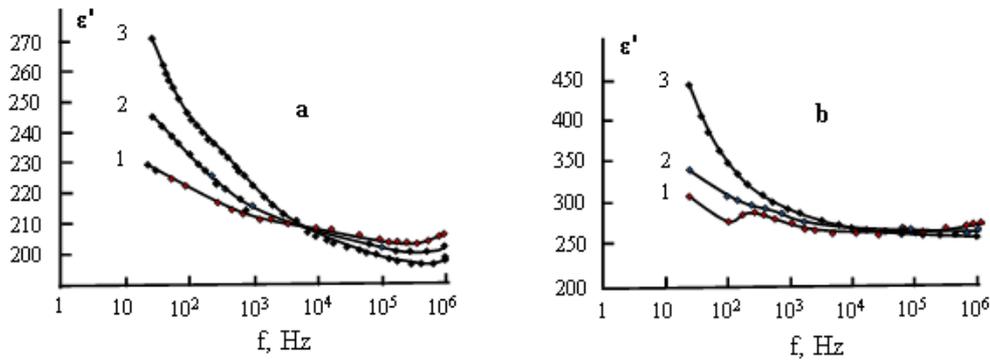


Fig. 6. Frequency dependences of the real part of the complex dielectric permittivity of $TiGa_{1-x}Er_xSe_2$ crystals: a) $x=0$; b) $x=0.005$. T: 1-180 K, 2-230 K, 3-300 K.

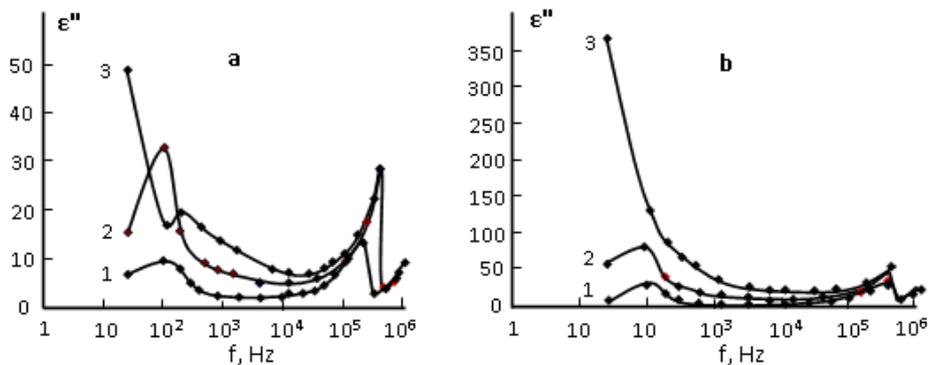


Fig. 7. Frequency dependences of the imaginary part of the complex dielectric permittivity of $TiGa_{1-x}Er_xSe_2$ crystals: a) $x=0$; b) $x=0.005$. T: 1-180 K, 2-230 K, 3-300 K.

We also investigated the frequency dependence of the ac-conductivity in TlGa_{1-x}Er_xSe₂ ($x=0; 0.005$) crystals (Fig. 9).

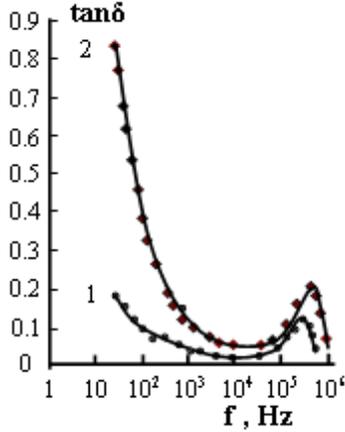


Fig. 8. Frequency dependences of dielectric loss tangent in TlGa_{1-x}Er_xSe₂ crystals: 1) $x=0$; 2) $x=0.005$, $T = 300$ K

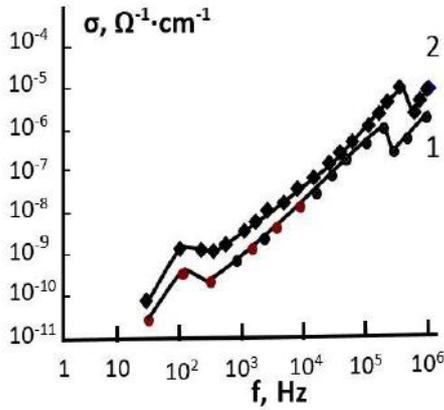


Fig. 9. Frequency dependence of the conductivity of TlGa_{1-x}Er_xSe₂ crystals: 1) $x = 0$, 2) $x=0.005$. $T=180$ K.

The σ_{ac} value for TlGa_{0.995}Er_{0.005}Se₂ solid solution is higher than for TlGaSe₂ crystal.

The dependence curves $\sigma_{ac}(f)$ of TlGa_{1-x}Er_xSe₂ ($x=0; 0.005$) solid solution are observed in 3 parts: in the first case, $\sigma_{ac} \sim f^{0.6}$ was observed and then increased to $\sigma_{ac} \sim f^{0.8}$ ($10^2 - 10^4$ Hz). Then, when the frequency was increased to 1 MHz, it was replaced by $\sigma_{ac} \sim f^{1.2}$. The fact that we obtain it from the law of $\sigma_{ac} \sim f^{0.8}$

indicates that there is the hopping conductivity in localized states near the Fermi level.

Table 4 shows the calculated parameters for TlGa_{1-x}Er_xSe₂ ($x=0; 0.005$) ($f = 10^4$ Hz) within the Mott approximation of solid solutions. The dependence of $\sigma_{ac} \sim f^{0.8}$ on the studied crystals at $10^2 - 10^4$ Hz frequencies is observed. The density of localized states (N_F) near the Fermi level for the TlGa_{1-x}Er_xSe₂ ($x=0; 0.005$) solid solutions is calculated according to Mott's theory [17] by the formula:

$$\sigma_{ac}(f) = \frac{\pi^3}{96} e^2 k T N_F^2 a^5 f \left[\ln \left(\frac{v_{ph}}{f} \right) \right]^4$$

Here e is the elementary charge, k - Boltzman constant, N_F - density of states near the Fermi level, $a = 1/a$ is the localization radius, $\psi \sim e^{-ar}$ is the wave function of the localized charge carriers; v_{ph} - phonon frequency. When calculating the density of localized states near the Fermi level, the localization radius $a = 34 \text{ \AA}$ was taken for the TlGaSe₂ crystal [15]. The value of the density of localized states (N_F) near the Fermi level calculated by Mott's theory is given in the table 1. When the temperature increases for 0.5 % Er samples, the value of the density of localized states near the Fermi level increases ($N_F = 2.23 \times 10^{18} \text{ eV}^{-1} \text{cm}^{-3}$ when $x = 0$, and $N_F = 3.08 \times 10^{18} \text{ eV}^{-1} \text{cm}^{-3}$ when $x=0.005$). According to the theory of hopping conductivity in the alternating current, the hopping distance (R) is determined by the expression [17]:

$$R = \frac{1}{2a} \ln \left(\frac{v_{ph}}{f} \right)$$

The value of the R and E_a was calculated for TlGa_{1-x}Er_xSe₂ crystals: $R=312 \text{ \AA}$ at 10^4 Hz [17]:

$$\Delta E = \frac{3}{2\pi R^3 N_F}, \quad E_a = \left(\frac{kT}{N_F a^3} \right)^{\frac{1}{4}}$$

According to this expression, the energy spread ΔE of the states localized near the Fermi level in TlGa_{1-x}Er_xSe₂ can be estimated.

Table 4 shows that the increase in temperature and the presence of Er in TlGaSe₂ increases the density of localized states near the Fermi level, narrows the energy spread ΔE , and changes the value of concentration on local levels.

Table 4.

The parameters of localized states in TlGa_{1-x}Er_xSe₂ solid solutions obtained from high-frequency dielectric measurements.

Crystal composition TlGa _{1-x} Er _x Se ₂	T, K	$N_F, \text{eV}^{-1} \text{cm}^{-3}$	$\Delta E, \text{meV}$	N_t, cm^{-3}	E_a, eV
$x = 0$	180	$1.54 \cdot 10^{18}$	10	$1.54 \cdot 10^{16}$	0.088
	230	$1.76 \cdot 10^{18}$	9	$1.56 \cdot 10^{16}$	0.102
	300	$2.23 \cdot 10^{18}$	7	$1.57 \cdot 10^{16}$	0.118
$x = 0.005$	180	$2.48 \cdot 10^{18}$	6.7	$1.66 \cdot 10^{16}$	0.078
	230	$2.58 \cdot 10^{18}$	6.3	$1.62 \cdot 10^{16}$	0.093
	300	$3.08 \cdot 10^{18}$	5	$1.54 \cdot 10^{16}$	0.109

The analysis of the absorption spectra of $\text{TlGa}_{1-x}\text{Er}_x\text{Se}_2$ crystals, namely, the dependence of $(\alpha\hbar\omega)^2$ on the photon energy $\hbar\omega$, made it possible to determine the energies of direct transitions in the studied crystals, and, consequently, their band gap E_g . In the temperature range from 77 to 300 K, the temperature dependences of the band gap of the studied solid solutions were constructed (Fig. 10).

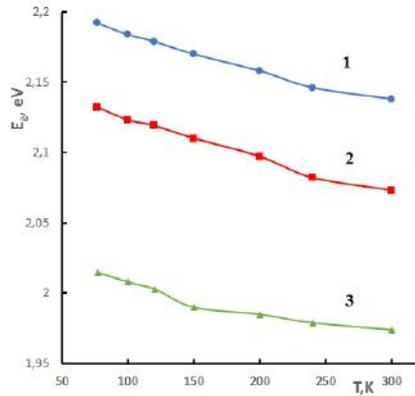


Fig. 10. Temperature dependences of the band gap of $\text{TlGa}_{1-x}\text{Er}_x\text{Se}_2$ single crystals. 1 - $x = 0$; 2 - $x = 0.001$; 3 - $x = 0.005$.

The experiment showed that the following differences are observed in the structure of the absorption edge of TlGaSe_2 single crystals and $\text{TlGa}_{1-x}\text{Er}_x\text{Se}_2$ solid solutions: at low temperatures, the

absorption band associated with the formation of an exciton near the direct edge cannot be detected in solid solutions; the absorption coefficient in $\text{TlGa}_{1-x}\text{Er}_x\text{Se}_2$ is noticeably higher than in TlGaSe_2 . The average temperature coefficient of the band gap dE_g/dT in the temperature range 77-300K for TlGaSe_2 and $\text{TlGa}_{0.999}\text{Er}_{0.001}\text{Se}_2$ is $-2.4 \cdot 10^{-4}$ eV/K, and for $\text{TlGa}_{0.995}\text{Er}_{0.005}\text{Se}_2$ is $-1.9 \cdot 10^{-4}$ eV/K. The long-wavelength band gap shift for $\text{TlGa}_{0.999}\text{Er}_{0.001}\text{Se}_2$ relative to TlGaSe_2 at $T=77$ K is 60 meV, and for $\text{TlGa}_{0.995}\text{Er}_{0.005}\text{Se}_2$ it is 177 meV.

4. CONCLUSIONS

The electrical conductivity, dielectric constant, and dielectric loss tangent of the $\text{TlGa}_{1-x}\text{Er}_x\text{Se}_2$ solid solution in an alternating electric field are investigated. In the crystals studied, the change in the dielectric constant as a function of temperature (150–300 K) at low frequencies is associated with the presence of relaxation polarization. $\sigma_{ac}(f)$ in $\text{TlGa}_{1-x}\text{Er}_x\text{Se}_2$ ($x = 0; 0.005$) crystals, at frequencies of 10^2 - 10^4 Hz obeys the law $\sigma_{ac} \sim f^{0.8}$. This indicates the presence of a hopping mechanism of charge transfer in states localized near the Fermi level. Doping the TlGaSe_2 crystal with erbium increases the density of states localized near the Fermi level and reduces their spread. An increase in the erbium concentration in $\text{TlGa}_{1-x}\text{Er}_x\text{Se}_2$ solid solutions leads to a decrease in the band gap and its temperature coefficient modulus.

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